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AMFR-W NUMERICAL METHODS FOR SOLVING HIGH DIMENSIONAL SABR/LIBOR PDE MODELS*

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Abstract. In this work we mainly develop a new numerical methodology to solve a PDE model 4 recently proposed in the literature for pricing interest rate derivatives. More precisely, we use high 5 6 order in time AMFR-W methods, which belong to a class of W-methods based on Approximate 7 Matrix Factorization (AMF) and are specially suitable in the presence of mixed spatial derivatives. High order convergence in time allows larger time steps which combined with the splitting of the 8 involved operators, highly reduces the computational time for a given accuracy. Moreover, the con-9 10 sideration of a large number of underlying forward rates makes the PDE problem high dimensional in space, so the use of AMFR-W methods with a sparse grids combination technique represents another 11 innovative aspect, making AMFR-W more efficient than with full grids and opening the possibility of 12 13 parallelization. Also the consideration of new homogeneous Neumann boundary conditions provides 14 another original feature to avoid the difficulties associated to the presence of boundary layers when using Dirichlet ones, specially in advection dominated regimes. These Neumann boundary condi-15 tions motivate the introduction of a modified combination technique to overcome a decrease in the 1617 accuracy of the standard combination technique.

18 Key words. SABR-LIBOR market models, high dimensional PDEs, AMFR-W methods, finite 19 differences, sparse grids combination technique

20 **AMS subject classifications.** 65M06, 65M20, 65M50, 65W10, 91G30, 91G80

21**1.** Introduction. High dimensional parabolic Partial Differential Equations (PDEs) arise in many fields of science and engineering problems, as for example in 22 computational biology for stochastic gene networks [2] or in computational finance for 23pricing financial derivatives [33], when a large number of underlying stochastic fac-24 tors are involved in their equivalent stochastic formulations. In fact, each stochastic 25factor gives rise to one spatial-like variable in the corresponding PDE. In this high 2627 dimensional setting, when using finite differences for the spatial discretization, the complexity of standard grid based approaches grows exponentially with the dimen-28 sions of the problem as well as the computational times, thus giving rise to the so 29called *curse of dimensionality*. Thus, alternative techniques to the standard full grid 30 are required. Also, the use of high order time integration schemes turns out to be very 31 convenient to allow larger time steps and therefore reduce the computational time to 33 get a prescribed accuracy.

In the present work we mainly propose a new numerical technique for solving the high dimensional PDE problem governing the Stochastic Alpha Beta Rho-LIBOR Market Model (SABR-LMM) PDE model introduced in [33]. The more classical LMM has been introduced for pricing interest rate derivatives which depend on the evolution

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of a certain number of forward LIBOR rates, when their volatility was assumed to 38 39 be constant (see [5], for example). More recently, the consideration of stochastic volatility has been incorporated for a better fit to market data by combining the 40 classical LMM and the SABR model for stochastic volatility in [21, 36, 38], with 41 different modelling approaches. In these works, the number of stochastic factors 42 depends on the number of forward rates and volatilities that are considered. Their 43 formulations are posed in terms of expectations to be solved by means of Monte Carlo 44 techniques. Although we are aware of the evolution of LMM as a consequence of the 45 financial crisis in 2007 to incorporate the practical presence of a multicurve setting 46(see [35], for example) and the recent ongoing studies related to LIBOR transitions 47 not consolidated in market practice yet [34], we have chosen to start from the classical 48 49 version of LMM as described in [5].

More recently, in [33] a PDE formulation is obtained for the Mercurio and Morini 50model presented in [36] and a set of numerical methods are proposed to solve it. More precisely, the combination of standard finite differences in space and a θ -method in 52time are proposed on uniform full grids. Moreover, by arguing that these standard 53 finite difference methods based on traditional full grids are not able to price interest 54rate derivatives with more than three or four stochastic forward rates, a sparse grid combination technique is applied. A rigorous analysis of finite differences schemes 56 in the sparse grid combination technique in arbitrary dimensions is carried out in 57 [39]. In order to overcome the curse of dimensionality one can try to use high order 58discretizations in time and space, mainly to reduce the required discretization points 60 to achieve certain accuracy. Although for sparse grids in space there are lots of works (see [26] and the references therein), for time discretization only schemes up to order 61 two have been applied. Following this objective, unlike in [33], in the present article 62 we propose higher order discretization techniques in time based on a special class of 63 W-methods [41], the AMFR-W methods introduced in [12]. 64

These AMFR-W-methods are specially suitable for parabolic problems involving 65 66 mixed spatial derivatives as it is the case in the SABR-LMM model proposed in [33]. In [12] it is shown that they are unconditionally stable regardless the spatial 67 dimension on linear constant coefficients PDEs with mixed derivative terms with 68 both periodic boundary conditions and homogeneous Dirichlet boundary conditions. 69 Moreover, as the ADI methods [26, 28], the AMFR-W scheme takes advantage of the 70 structure of the linear system obtained from the spatial discretization, so it can be 7172 decomposed into tridiagonal systems which can be solved in linear run-time. Thus, the computational effort is significantly reduced. However, while classical ADI schemes 73 are of order two in time, the proposed AMFR-W scheme exhibits order three in time 74 when standard full grids are used. An additional innovative aspect of the present work 76 is the application of the AMFR-W methods in the context of sparse grids combination technique, which turns out to be an efficient tool for solving the SABR-LMM model 77 in the required high dimensional setting. Furthermore, the introduction of more 78 appropriate Neumann boundary conditions motivates the consideration of a modified 79 combination technique to improve the convergence. 80

LMMs are usually simulated by means of Monte Carlo method, in contrast with the here proposed PDE methodology. We aim to avoid the slow Monte Carlo rate of convergence, $O(1/\sqrt{M})$ for all dimensions, M being the number of simulations.

The plan of the article is the following. In Section 2 we present the PDE model and justify the introduction of new homogeneous Neumann boundary conditions at the so called outflow boundaries. In Section 3, we introduce the space discretization of the PDE problem with finite differences to obtain an ODE system in suitable form for the application of the AMFR-W method. In Section 4 we describe the application of AMFR-W method to the ODE system to get the fully discretized problem. Section 5 is devoted to the methodology of sparse grids, including the standard and the modified combination techniques. In Section 6 we present and discuss the numerical results

92 obtained for full and sparse grids. Finally, Section 7 contains some conclusions.

2. PDE formulation of the SABR-LMM model. As indicated in the pre-93 vious section, we mainly address a new and more efficient numerical solution of the 94 PDE formulation introduced in [33] for the SABR-LMM proposed by Mercurio and 95 Morini in [36] to price a financial derivative which depends on a certain number of 96 97 forward LIBOR rates, i.e a swaption. In this section we introduce the PDE model and we incorporate some new boundary conditions. Although we refer the reader to 98 99 [33] for the statement of the model and further details, we need to introduce some financial concepts and their notations related to interest rates derivatives to be used 100 along this article. In this respect, we also refer the reader to the textbook [5]. 101

A zero coupon *bond* with maturity at time T pays its holder one unit of currency at time T. The zero coupon value at time t < T will be denoted by P(t,T), and is also referred as the discount factor from time T to time t. A *tenor structure* is defined as a set of ordered payment dates $T_0 < T_1 < \ldots < T_{N-1} < T_N$. The time gap between two consecutive tenor dates is denoted by $\tau_i = T_{i+1} - T_i$. In view of previous definitions, a payment of x units at time T_i is worth $xP(t,T_i)$ at time $t < T_i$.

Next, we introduce the forward rates that enter in the LMM. We consider the forward interest rate $F_i(t)$ as an interest rate we can contract to borrow or lend money during the future time period $[T_i, T_{i+1}]$, which is fixed at time T_i . Moreover, the value of $F_i(t)$ can be expressed in terms of discount factors in the form:

$$F_{i}(t) = F(t; T_{i}, T_{i+1}) = \frac{1}{\tau_{i}} \left(\frac{P(t, T_{i})}{P(t, T_{i+1})} - 1 \right) \text{ where } t \leq T_{i}$$

Conversely, the price of a zero coupon bond at time T_i that matures at T_j , $P(T_i, T_j)$, can be expressed in terms of forward LIBOR rates as

$$P(T_i, T_j) = \prod_{k=i}^{j-1} \frac{1}{1 + \tau_k F_k(T_i)}$$

Among all interest rate derivatives, the simplest one is the caplet. A *caplet* is a European call option on a forward rate. Thus, if the maturity of a caplet is T_{i+1} , at that time the holder of the caplet receives the payoff $\tau_i(F_i(T_i) - K)^+$, so its discounted payoff at time $t < T_{i+1}$ is given by $P(t, T_{i+1})\tau_i(F_i(T_i) - K)^+$, where $(\cdot)^+$ denotes the function max $(\cdot, 0)$ and K is the strike (a fixed interest rate) of the caplet. If constant volatilities are assumed as it is the case in the classical LMM, the caplet price can be analytically computed with a Black's formula (see [5], for details).

115 An interest rate *swap* (IRS) is a contract to exchange interest payments at future 116 fixed dates. At every time instant in the prescribed set of dates T_{a+1}, \ldots, T_b the 117 contract holder pays a fixed interest rate K and receives a floating forward LIBOR 118 rate $F_i(T_i)$, which is fixed at time T_i . At time T_a the value of the IRS is given by

119 (2.1)
$$\operatorname{IRS}(T_a; T_a, \dots, T_b) = \sum_{i=a}^{b-1} P(T_a, T_{i+1}) \tau_i(F_i(T_a) - K).$$

120 A European $T_a \times (T_b - T_a)$ swaption is an option that gives the right to enter a 121 swap at the future time T_a (swaption maturity). The underlying swap length $T_b - T_a$ is referred as the tenor of the swaption. Therefore, the discounted swaption payoff to time t is equal to $P(t, T_a)(\text{IRS}(T_a; T_a, \dots, T_b))^+$.

In the forthcoming section devoted to numerical results, several examples address the pricing of caplets and swaptions. Note that the payoff of a caplet just involves one forward rate, therefore its price at any time before maturity only depends on this particular forward rate. In the case of IRS or swaptions, their payoffs at expiration date depend on a certain number of forward rates, so their price at any time before expiration will also depend on them.

In the Mercurio and Morini model it is assumed that a generic European interest 130rate derivative depends on the evolution of N-1 forward rates, $F_1, F_2, \ldots, F_{N-1}$ 131associated to bonds with maturity related to the tenor structure, and a common 132stochastic volatility, V. Let $u = u(t, F_1, F_2, \ldots, F_{N-1}, V)$ denote the value at time $t \in$ 133[0,T] of this European interest rate derivative, with forward rates $F_i \in [0, F_i^{max}], i =$ 134 $1, 2, \ldots, N-1$, and volatility $V \in [0, V^{max}]$. Note that the previously described 135caplets, IRS and swaptions are particular cases. As stated in [33], the function u136137satisfies the following PDE

$$\begin{split} \frac{\partial u}{\partial t} &+ \frac{1}{2} \sigma^2 V^2 \frac{\partial^2 u}{\partial V^2} + \frac{1}{2} V^2 \sum_{i,j=1}^{N-1} \alpha_i \alpha_j \rho_{ij} F_i^\beta F_j^\beta \frac{\partial^2 u}{\partial F_i \partial F_j} \\ &+ \sigma V^2 \sum_{i=1}^{N-1} \alpha_i \phi_i F_i^\beta \frac{\partial^2 u}{\partial F_i \partial V} + \sum_{i=1}^{N-1} \mu_i F_i^\beta \frac{\partial u}{\partial F_i} \quad = 0, \end{split}$$

where σ is the volatility of the stochastic volatility V, α_i is a deterministic (constant) 139instantaneous volatility coefficient of the forward rate F_i , ρ_{ij} is the correlation between 140the forward rates F_i and F_j , ϕ_i is the correlation between F_i and the stochastic 141 volatility V and μ_i is the drift of the *i*-th forward rate. Moreover, when the bond 142 $P(t,T_1)$ is chosen as the numeraire, the drifts μ_i depend on the forward rates as 143 $\mu_1 = 0, \ \mu_i = \alpha_i V^2 \sum_{j=2}^i \frac{\tau_j F_j^{\beta}}{1 + \tau_j F_j} \rho_{ij} \alpha_j, \ i \ge 2.$ The parameter $\beta \in [0, 1]$ is the elasticity 144of variance, which usually is 0, 0.5 or 1, that corresponds to stochastic volatilities 145 with normal, CIR or log-normal dynamics. For the correlation structure, as in [33] we consider the expression $\rho_{ij} = e^{-\lambda |T_i - T_j|}$, which depends on the constant parameter λ . 146147In view of the form of the differential operator governing the PDE, (2.2) must be 148completed with a final condition $u(T, F_1, F_2, ..., F_{N-1}, V) = q(T, F_1, F_2, ..., F_{N-1})$ 149where q represents the derivative payoff, the expression of which depends on the 150151interest rate derivative we are dealing with. In order to apply the method of lines (MoL) to discretize the previous model

In order to apply the method of lines (MoL) to discretize the previous model it is more convenient to write (2.2) in terms of the time to maturity T - t instead of the physical time t, so that the final condition turns into an initial condition. In an abuse of notation we keep the notation t for the new formulation after this change in the time variable. More precisely, we rewrite the model by denoting u = $u(t, F_1, F_2, \ldots, F_{N-1}, F_N)$ the value of the interest rate derivative at time $T - t, t \in$ [0, T], with forward rates $F_i \in [0, F_i^{max}], i = 1, 2, \ldots, N - 1$, and volatility $V \in$ $[0, V^{max}]$. After some easy calculus, we obtain the equation

160 (2.3)
$$\frac{\partial u}{\partial t} = \sum_{i=1}^{N} d_i \frac{\partial^2 u}{\partial F_i^2} + \sum_{i=1}^{N-1} \sum_{k=i+1}^{N} m_{ik} \frac{\partial^2 u}{\partial F_i \partial F_k} + \sum_{i=2}^{N-1} a_i \frac{\partial u}{\partial F_i},$$

161 with the *initial condition*

162 (2.4)
$$u(0, F_1, F_2, \dots, F_{N-1}, F_N) = g(T, F_1, F_2, \dots, F_{N-1}),$$

163 where

175

164 (2.5)
$$d_i = d_i(F_i, F_N) = \begin{cases} \frac{1}{2} \alpha_i^2 \rho_{ii} F_i^{2\beta} F_N^2, & \text{if } 1 \le i \le N-1, \\ \frac{1}{2} \sigma^2 F_N^2, & \text{if } i = N. \end{cases}$$

165 Moreover, for i = 1, ..., N - 1, k = i + 1, ..., N, we define

166 (2.6)
$$m_{ik} = m_{ik}(F_i, F_k, F_N) = \begin{cases} \alpha_i \alpha_k \rho_{ik} F_i^{\beta} F_k^{\beta} F_N^2, & \text{if } i+1 \le k \le N-1, \\ \alpha_i \sigma \phi_i F_i^{\beta} F_N^2, & \text{if } k = N. \end{cases}$$

167 Note that the last term in (2.3) is only defined for $N \ge 3$. In this case, for each 168 i = 2, ..., N - 1 we define:

169 (2.7)
$$a_i = a_i(F_2, \dots, F_i, F_N) = \left(\sum_{j=2}^i \alpha_i \alpha_j \rho_{ij} \Phi_\beta(F_j, \tau_j)\right) F_i^\beta F_N^2,$$

170 where Φ_{β} is the scalar function $\Phi_{\beta}(x,\tau) := \frac{\tau x^{\beta}}{1+\tau x}, \tau > 0, x \ge 0.$

In next paragraphs we will discuss about the appropriate boundary conditions to add to (2.3)-(2.4) to define the initial-boundary value problem. First, note that in [33] the following time-independent boundary conditions were considered to complete the formulation (2.3)-(2.4):

(2.8)
if
$$F_j = 0$$
 or $F_j = F_j^{max}$, $1 \le j \le N - 1$,
 $u(t, F_1, \dots, F_{N-1}, F_N) = g(T, F_1, \dots, F_{N-1})$,
 $u(t, F_1, \dots, F_{N-1}, 0) = u(0, F_1, \dots, F_{N-1}, 0) = g(T, F_1, \dots, F_{N-1})$,
 $\frac{\partial u}{\partial V}(t, F_1, \dots, F_{N-1}, V^{max}) = 0$.

These boundary conditions are appropriate when $\beta = 0$ since the PDE coefficients 176are independent of the forward rates $(F_j, j \leq N-1)$ and the advection terms are 177 moderate. When $\beta > 0$, they also can be appropriate when we have a small number 178of forward rates and a derivative with short maturity. Nevertheless, as soon as this 179number increases then the dimension of the PDE grows, and the advection-dominance 180of the PDE (2.3) becomes more relevant. It is well-known that imposing a Dirichlet 181 boundary condition at the outflow boundaries $F_i = F_i^{max}$, $i = 1, \ldots, N-1$, in an 182advection-dominated setting could give rise to boundary layers (see, for instance, [27, 183 Sect. I.5]), and some previous numerical results with this model corroborate that. 184

185 When $\beta \in (0, 1]$, in order to show that the larger the dimension N the more 186 advection-dominant the PDE (2.3) becomes, it is better to write the PDE (2.3) in the 187 following conservative form

188 (2.9)
$$u_t + \nabla \cdot (\underline{c}u) = \nabla \cdot (D\nabla u) + s(\underline{x}, t, u), \quad \underline{x} \in \Omega \subset \mathbb{R}^N, t > 0,$$

189 where $\underline{x} = (F_1, \dots, F_N), \ \underline{c} = (c_1, \dots, c_N)^T, \ D = (D_{ik})_{i,k=1}^N, \ c_i = c_i(\underline{x}, t), \ D_{ik} = 190 \quad D_{ik}(\underline{x}, t).$

191 Note that in multi-dimensional advection-diffusion-reaction PDEs with variable 192 coefficients, some additional advection terms arise in the conservative form coming 193 from the partial derivatives of the coefficients of the second order diffusion terms. 194 After some manipulations it can be proved that PDE (2.3) admits an expression 195 of type (2.9) with $D_{ii} = d_i$, $D_{ik} = m_{ik}/2$, $\forall i \neq k$, and

(2.10) $c_i = \left(P_i^+ - P_i^-\right) \alpha_i F_i^\beta F_N^2 + \alpha_i \phi_i F_i^\beta \sigma F_N, \quad 1 \le i \le N - 1, \\ c_N = P_N^+ \sigma F_N^2 + \sigma^2 F_N,$

where

$$P_{i}^{+} = \alpha_{i} \frac{\beta}{F_{i}^{1-\beta}} + \sum_{\substack{j=1\\j\neq i}}^{N-1} \frac{1}{2} \alpha_{j} \rho_{ij} \frac{\beta}{F_{j}^{1-\beta}}, \quad 1 \le i \le N-1,$$

$$P_{N}^{+} = \sum_{j=1}^{N-1} \frac{1}{2} \alpha_{j} \phi_{j} \frac{\beta}{F_{j}^{1-\beta}}, \quad P_{1}^{-} = 0, \quad P_{2}^{-} = \alpha_{2} \Phi_{\beta}(F_{2}, \tau_{2}),$$

$$P_{i}^{-} = \alpha_{i} \Phi_{\beta}(F_{i}, \tau_{i}) + \sum_{j=2}^{i-1} \alpha_{j} \rho_{ij} \Phi_{\beta}(F_{j}, \tau_{j}), \quad 3 \le i \le N-1.$$

197 Moreover, the reaction term takes the form $s(\underline{x}, t, u) = \delta u$, $\delta = P F_N^2 + Q \sigma F_N + \sigma^2$, 198 where P and Q depend on the forward rates F_1, \ldots, F_{N-1} .

199 Obviously, when $0 < \beta < 1$ the PDE (2.3) is equivalent to (2.9) whenever $F_j >$ 200 $0, \forall j = 1, ..., N - 1$, due to the lack of differentiability of \underline{c} when $F_j = 0$. Since 201 we intend to use this expression (2.9) to study the behaviour of the PDE at the 202 boundaries $F_j = F_j^{max}$, we will assume for now that $F_j > 0, 1 \le j \le N - 1$. On the 203 other hand, when $\beta = 1$ this additional assumption is not needed.

In most of practical cases we have that

what implies that, for $2 \le i \le N - 1^1$, (2.12)

$$P_{i}^{+} - P_{i}^{-} = \alpha_{i} \left(\frac{\beta}{F_{i}^{1-\beta}} - \Phi_{\beta}(F_{i}, \tau_{i}) \right) + \sum_{j=2}^{i-1} \frac{1}{2} \alpha_{j} \rho_{ij} \left(\frac{\beta}{F_{j}^{1-\beta}} - 2\Phi_{\beta}(F_{i}, \tau_{i}) \right) + \frac{1}{2} \alpha_{1} \rho_{i1} \frac{\beta}{F_{1}^{1-\beta}} + \sum_{j=i+1}^{N-1} \frac{1}{2} \alpha_{j} \rho_{ij} \frac{\beta}{F_{j}^{1-\beta}} \ge 0.$$

We must observe that in the case $\beta = 1$,

$$\frac{\beta}{F_i^{1-\beta}} - \Phi_{\beta}(F_i, \tau_i) = 1 - \frac{\tau_i F_i}{1 + \tau_i F_i} = \frac{1}{1 + \tau_i F_i} \ge 0,$$
$$\frac{\beta}{F_j^{1-\beta}} - 2\Phi_{\beta}(F_j, \tau_j) = 1 - 2\frac{\tau_i F_i}{1 + \tau_i F_i} = \frac{1 - \tau F_j}{1 + \tau_i F_j} \ge 0 \quad \text{if } \tau_j F_j \le 1,$$

and the second requirement in (2.11) is superfluous.

Therefore, $c_i \ge 0$, i = 1, ..., N, so it is clear that all the boundaries $F_i = F_i^{max}$ 10 $1 \le i \le N$ are outflow boundaries since the outward normal vector on each one of these

borders is the *N*-dimensional canonical vector $\mathbf{e}_i = (e_{ik})_{k=1}^N$, $e_{ii} = 1$, $e_{ik} \neq 0$, $\forall k \neq i$.

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¹In the sequel, $\sum_{i=j}^{k} (\cdot) = 0$ when j > k.

On the other hand, we need to take into account that these advection coefficients depend on N, $c_i = c_i^{(N)}$, and $F_N = V$ so if we increase the number of forward rates from $\{F_1, \ldots, F_{N-2}\}$ to $\{F_1, \ldots, F_{N-2}, F_{N-1}\}$, we have that

215 (2.13)
$$c_i^{(N)} = c_i^{(N-1)} + \frac{1}{2}\alpha_i\alpha_{N-1}\rho_{i,N-1}\frac{\beta}{F_{N-1}^{1-\beta}}F_i^{\beta}V^2, \quad i = 1, \dots, N-2.$$

Since usually $F_i^{max} = F^{max}$, i = 1, ..., N - 1 and $V^{max} >> F^{max}$, we can see that the advection increases with N whereas the diffusion coefficients do not.

A way to avoid the boundary layers that this advection-dominance can produce with conditions (2.8) when $0 < \beta \leq 1$, is to consider homogeneous Neumann boundary conditions instead [27], i.e.,

221

$$u(t, F_1, \dots, F_{N-1}, F_N) = g(T, F_1, \dots, F_{N-1}), \quad \text{if } F_j = 0, \ 1 \le j \le N,$$
$$\frac{\partial u}{\partial F_j}(t, F_1, \dots, F_{N-1}, F_N) = 0, \quad \text{if } F_j = F_j^{max}, \ 1 \le j \le N-1, \ \text{or } F_N = V^{max}.$$

On the one hand, these homogeneous Neumann boundary conditions are appropriate from the financial point of view. Actually, if we analyse the behaviour of the payoff of the swaption $T_1 \times (T_N - T_1)$, which is given by

$$g(T, F_1, F_2, \dots, F_{N-1}) = \max\left\{\sum_{i=1}^{N-1} \frac{\tau_i(F_i - K)}{(1 + \tau_1 F_1) \cdots (1 + \tau_i F_i)}, 0\right\},\$$

when some forward rate $F_j \to \infty$, we can consider the approximation

$$g(T, F_1, \dots, F_{N-1}, V) \approx f(F_1, \dots, F_{N-1}) := \sum_{i=1}^{N-1} \frac{\tau_i(F_i - K)}{P_i}$$

222 where $P_i := \prod_{l=1}^{i} (1 + \tau_l F_l), 1 \le i \le N - 1$, and prove that

223 (2.15)
$$\frac{\partial f}{\partial F_j}(T, F_1, \dots, F_{N-1}, V) \longrightarrow 0 \text{ when } F_j \to \infty, \quad j = 1, \dots, N-1.$$

In order to prove this, it is enough to write the partial derivatives of f as

225
$$\frac{\partial f}{\partial F_j}(F_1,\ldots,F_{N-1}) =$$

226
$$\frac{\tau_j}{(1+\tau_1F_1)\cdots(1+\tau_{j-1}F_{j-1})(1+\tau_jF_j)^2}\left((1+\tau_jK)-\sum_{i=j+1}^{N-1}\frac{\tau_i(F_i-K)}{\tilde{P}_{ij}}\right),$$

where $\tilde{P}_{ij} = P_i/P_j = (1+\tau_{j+1}F_{j+1})\cdots(1+\tau_iF_i), i \ge j+1$. As in [33], an homogeneous Neumann boundary condition is imposed at $F_N = V^{max}$ because the price of the derivative becomes independent of V when V approaches to infinity.

On the other hand, the choice (2.14) becomes better when some discretization of the "spatial" variables F_j is used to approximate the solution of PDE (2.9) on a uniform spatial grid, as will be illustrated later on. **3.** Space discretization with finite differences. Following the ideas given in [33], firstly a space discretization is performed on a uniform spatial grid on $\Omega =$ $[0, F_1^{max}] \times [0, F_2^{max}] \times \cdots \times [0, F_{N-1}^{max}] \times [0, V^{max}]$. However, we propose here a different time discretization by using the AMFR-W-methods introduced in [12].

For the space discretization we consider N integers $(M_1, \ldots, M_{N-1}, M_N)$ to define on Ω the spatial grid with $M_i + 1$ equally spaced points (denoting $F_N^{max} = V^{max}$) at the F_i -direction $F_{i,j_i} = j_i h_i$, $0 \le j_i \le M_i$, $h_i = \frac{F_i^{max}}{M_i}$, $1 \le i \le N$, and discretize the derivatives in (2.3) with second order central finite differences at each spatial node $(F_{1,j_1}, \ldots, F_{N-1,j_{N-1}}, F_{N,j_N})$. More precisely, the MoL approximates the solution at each spatial point $u(t, F_{1,j_1}, \ldots, F_{N-1,j_{N-1}}, F_{N,j_N}) \approx U_{j_1,\ldots,j_{N-1},j_N}(t)$, where the values $U_{j_1,\ldots,j_{N-1},j_N}(t)$ need to satisfy the semi-discretized ODE system

$$\frac{d}{dt}U_{j_1,\dots,j_{N-1},j_N}(t) = \sum_{\substack{i=1\\N-1}}^{N} d_i(F_{i,j_i},F_{N,j_N}) \Delta_{j_1,\dots,j_{N-1},j_N}^{(i)} + \sum_{\substack{i=1\\N-1}}^{N} \sum_{\substack{k=i+1\\k=i+1}}^{N} m_{ik}(F_{i,j_i},F_{k,j_k},F_{N,j_N}) \Delta_{j_1,\dots,j_{N-1},j_N}^{(ik)} + \sum_{\substack{i=2\\i=2}}^{N} a_i(F_{2,j_2},\dots,F_{i,j_i},F_{N,j_N}) \nabla_{j_1,\dots,j_{N-1},j_N}^{(i)},$$

with $\Delta^{(i)}$, $\Delta^{(ik)}$, $\nabla^{(i)}$ representing the approximations with central differences of the derivatives $(\partial^2 u/\partial F_i^2)$, $(\partial^2 u/\partial F_i \partial F_k)$, $(\partial u/\partial F_i)$, respectively.

One important drawback of finite differences comes from the complexity of dealing with all these approximations when the dimension N is large. For the efficient manipulation of these differences, we propose to use N-dimensional multi-indices $\mathbf{j} = (j_1, \ldots, j_N)$, and the following Lemma 3.1 that is proved in the Appendix.

LEMMA 3.1. Given N pairs of integers $m_i \leq M_i$ for i = 1, ..., N, let us define the set $\mathcal{I}_N = \{\mathbf{j} = (j_1, ..., j_N) \mid m_i \leq j_i \leq M_i, \forall i = 1, ..., N\}$. Moreover, for $M_T = \prod_{k=1}^N (M_k - m_k + 1)$, define the map $\vartheta : \mathcal{I}_N \longrightarrow \{m_1, m_1 + 1, ..., M_T + m_1 - 1\}$,

255 (3.2)
$$\vartheta(\mathbf{j}) = J = j_1 + \sum_{l=2}^{N} \left((j_l - m_l) \prod_{r=1}^{l-1} (M_r - m_r + 1) \right).$$

- 256 Then, the map ϑ is bijective.
- Besides the proof of Lemma 3.1, in the Appendix it is also included a practical way for computing the inverse map $\vartheta^{-1}(J) = \mathbf{j}$ (see (A.3) in Lemma A.1).

Note that due to the Dirichlet boundary conditions in (2.14), the values for $j_i = 0$ for $1 \leq i \leq N$ are given by the derivative payoff. As a consequence, the ODE system (3.1) is applied only when $j_i = 1, \ldots, M_i, 1 \leq i \leq N$, so it has dimension $L = M_1 \cdots M_{N-1} M_N$. Then, we separate the multi-indices that correspond with finite differences nodes on the lower boundaries from the rest of them, so that two different bijections of type (3.2) are considered:

$$\vartheta_0 : \mathcal{I}_N^{(0)} = \{ \mathbf{j} = (j_1, \dots, j_N) \mid 0 \le j_i \le M_i, \ \forall i = 1, \dots, N \} \longrightarrow \{0, 1, \dots, M - 1\}, \\ \vartheta_1 : \mathcal{I}_N^{(1)} = \{ \mathbf{k} = (k_1, \dots, k_N) \mid 1 \le k_i \le M_i, \ \forall i = 1, \dots, N \} \longrightarrow \{1, 2, \dots, L\},$$

where $M = \prod_{k=1}^{N} (M_k + 1)$, $\mathcal{I}_N^{(1)} \subset \mathcal{I}_N^{(0)}$, and $\vartheta_0(\mathbf{j}) = J = j_1 + \sum_{l=1}^{N} \left(j_l \prod_{l=1}^{l-1} (M_r + 1) \right)$,

$$\vartheta_{1}(\mathbf{k}) = S = J_{1} + \sum_{l=2}^{N} \left(J_{l} \prod_{r=1}^{n} (M_{r} + 1) \right), \quad \mathbf{j} \in \mathcal{I}_{N}^{n},$$
$$\vartheta_{1}(\mathbf{k}) = K = k_{1} + \sum_{l=2}^{N} \left((k_{l} - 1) \prod_{r=1}^{l-1} M_{r} \right), \quad \mathbf{k} \in \mathcal{I}_{N}^{(1)}.$$

 $\mathbf{i} \in \tau^{(0)}$

The set of integers $\{0, 1, \ldots, M-1\}$ can be obtained as the union of the two disjoint sets $\mathsf{Inner} = \vartheta_0 \left(\mathcal{I}_N^{(1)} \right) = \vartheta_0 \left(\vartheta_1^{-1} \left(\{1, \ldots, L\} \right) \right)$, $\mathsf{Outer} = \{0, 1, \ldots, M-1\} - \mathsf{Inner}$, and we consider the vector $Y(t) = (Y_J(t))_{J=0}^{M-1}$, where for each $J = 0, \ldots, M-1$, $(j_1, \ldots, j_{N-1}, j_N) = \vartheta_0^{-1}(J)$,

263 (3.3)
$$Y_J(t) = \begin{cases} U_{j_1,\dots,j_{N-1},j_N}(t), & \text{if } J \in \mathsf{Inner}, \\ g(T,F_{1,j_1},\dots,F_{N-1,j_{N-1}}), & \text{if } J \in \mathsf{Outer}. \end{cases}$$

Note that $\vartheta_0(\mathbf{j} - \mathbf{e}_i) \in \{0, 1, \dots, M - 1\}$, for all $J \in \mathsf{Inner}$ with $\mathbf{j} = \vartheta_0^{-1}(J) \in \mathcal{I}_N^{(1)}$. Moreover, $\vartheta_0(\mathbf{j} - \mathbf{e}_i) = J - E_i$, for $E_1 = 1$, $E_i = \prod_{r=1}^{i-1}(M_r + 1)$, $i \ge 2$. On the other hand, taking into account boundary conditions (2.14), for all $J \in \mathsf{Inner}$ we have that $\mathbf{j} = \vartheta_0^{-1}(J) \in \mathcal{I}_N^{(1)}$ and $\vartheta_0(\mathbf{j} + \mathbf{e}_i) \in \{0, 1, \dots, M - 1\}$ except when $j_i = M_i$, since $\mathbf{j} + \mathbf{e}_i = (\dots, M_i + 1, \dots) \notin \mathcal{I}_N^{(0)}$. So, if $J \in \mathsf{Inner}$, then $J + E_i = \vartheta_0(\mathbf{j} + \mathbf{e}_i)$, for all $i = 1, \dots, N$. However, for the case $j_i = M_i$ we take the virtual value $Y_{J+E_i} = Y_{J-E_i}$ due to the homogeneous Neumann conditions (2.14) at these boundaries. Therefore, taking into account (3.3), the ODE system (3.1) is shown in Sketch 1.

286 end for

10

287
$$Y'_{J}(t) = \sum_{i=1}^{N} (d_{i})_{J} \Delta_{J}^{(i)} + \sum_{i=1}^{N-1} \sum_{k=i+1}^{N} (m_{ik})_{J} \Delta_{J}^{(ik)} + \sum_{i=2}^{N-1} (a_{i})_{J} \nabla_{J}^{(i)}$$

end for 288

Thus, we obtain the semi-discretized autonomous initial value problem (IVP) 289

290 (3.4)
$$Y' = \mathcal{F}(Y), \ Y(0) = Y_0, \ t \in [0,T], \qquad \mathcal{F}(Y) = \mathcal{F}_0(Y) + \sum_{i=1}^N \mathcal{F}_i(Y).$$

In the splitting of the derivative \mathcal{F} , for each $i = 1, \ldots, N$, the term $\mathcal{F}_i(Y)$ contains 291 292the second order differences in the F_i -direction. That is, for all $J = 0, 1, \ldots, M - 1$,

293 (3.5)
$$(\mathcal{F}_{i}(Y))_{J} = \begin{cases} (d_{i})_{J} \Delta_{J}^{(i)}, & \text{if } J \in \text{Inner}, (j_{1}, \dots, j_{N}) = \vartheta_{0}^{-1}(J), \\ 0, & \text{if } J \in \text{Outer}, \end{cases}$$

while the term $\mathcal{F}_0(Y)$ gathers the discretization corresponding to the remaining terms. 294Clearly, $\forall i = 0, 1, \dots, N, \mathcal{F}_i(Y) = \mathcal{A}_i Y$, where for all $J \in \mathsf{Outer}$, the J-th row of 295

the constant matrix \mathcal{A}_i is null. Besides, when $i \geq 1$ and $J \in \mathsf{Inner}$, the J-th row of \mathcal{A}_i 296 has at most three non-zero elements, located at the columns $J - E_i$, J and $J + E_i$. 297Therefore, the differential system $Y' = \mathcal{F}(Y)$ in (3.4) involves actually L unknowns 298 since the corresponding equation for each $J \in \mathsf{Outer}$ is null. Once the semi-discretized 299linear IVP (3.4) has been posed, the AMFR-W-methods given in [12] can be applied 300 for its time integration. 301

4. Time discretization. 302

4.1. AMFR-W-methods. AMFR-W-methods belong to the class of W-meth-303 304 ods [41] for the time integration of IVPs of type (3.4). Thus, from an approximation Y_n of the solution Y(t) at $t = t_n$ and a step size $\Delta t > 0$, an s-stage W-method gives 305 the approximation Y_{n+1} at $t_{n+1} = t_n + \Delta t$ by 306

(4.1)
$$(I - \theta \Delta t W) K_r = \Delta t \mathcal{F} \Big(Y_n + \sum_{j=1}^{r-1} a_{rj} K_j \Big) + \sum_{j=1}^{r-1} q_{rj} K_j, \quad r = 1, 2, \dots, s,$$
$$Y_{n+1} = Y_n + \sum_{r=1}^s b_r K_r.$$

Each W-method is characterized by its coefficients (A, Q, b, θ) , where $A = (a_{rj})_{j < r}$, 308 $Q = (q_{rj})_{j < r}$ and $b = (b_r)$, and by the arbitrary matrix W. These kind of methods 309 can also be understood as a generalization of Rosenbrock methods, which are obtained 310 when $W = \mathcal{F}'(Y_n)$. In order to get W-methods of high order, W must be some rough 311 approximation of $\mathcal{F}'(Y_n)$, and methods of order 3 and higher can be found in the 312 313 literature under the assumption (see, for instance, [11, 37, 29, 15])

314 (4.2)
$$W - \mathcal{F}'(Y_n) = \mathcal{O}(\Delta t), \qquad \Delta t \to 0.$$

When $\mathcal{F}'(Y)$ admits a directional splitting of type (3.4), i.e. $\mathcal{F}'(Y_n) = \mathcal{F}'_0(Y_n) + \mathcal{F}'_0(Y_n)$ 315 $\sum_{i=1}^{N} \mathcal{A}_i$, where the matrices \mathcal{A}_i have simple structures, in [12] the authors propose 316(4.3)

317
$$\left(I - \theta \Delta t W\right)^{-1} = \prod_{i=N}^{1} \left(I - \nu \Delta t \mathcal{A}_i\right)^{-1} \left(2I - \left(I - \theta \Delta t \mathcal{F}'(Y_n)\right) \prod_{i=N}^{1} \left(I - \nu \Delta t \mathcal{A}_i\right)^{-1}\right),$$

resulting AMFR-W-method (4.1)-(4.3) is computed as

$$K_{r}^{(0)} = \Delta t \mathcal{F}(Y_{n} + \sum_{j=1}^{r-1} a_{rj}K_{j}) + \sum_{j=1}^{r-1} q_{rj}K_{j},$$

$$(I - \nu\Delta t\mathcal{A}_{i})K_{r}^{(i)} = K_{r}^{(i-1)}, \quad i = 1, \dots, N,$$

$$\hat{K}_{r}^{(0)} = 2K_{r}^{(0)} - (I - \theta\Delta t\mathcal{F}'(Y_{n}))K_{r}^{(N)},$$

$$(I - \nu\Delta t\mathcal{A}_{i})\hat{K}_{r}^{(i)} = \hat{K}_{r}^{(i-1)}, \quad i = 1, \dots, N,$$

$$K_{r} = \hat{K}_{r}^{(N)}.$$

In [12], different choices for the coefficients of these methods are tested. In this article we have used the proposed **AMFR-W2**, that is a 2-stage AMFR-W-method, with coefficients [27, p. 400]

328 (4.5)
$$a_{21} = 2/3, \quad q_{21} = -4/3, \quad b_1 = 5/4, \quad b_2 = 3/4.$$

Since (4.2) is fulfilled, this method is of order 3 for $\theta = (3 + \sqrt{3})/6$ (in ODE sense). In 329 330 [12], the authors also proved that the choice of the parameter ν depends on the number N of terms in the splitting in (3.4) to get unconditional stability. More precisely, they 331 apply this method on the parabolic test problem given in [13] and guarantee that 332 this method is unconditionally stable on linear constant coefficients PDEs with mixed 333 derivatives of dimension N if $\nu \geq N\kappa_N \theta$ with the values of κ_N given in [12, Table 334 335 2], when both periodic and homogeneous Dirichlet boundary conditions. In [16], the authors solve the two-dimensional PDE for the well-known Heston model in options 336 pricing. For this purpose, a hyperbolic change of variables is previously applied to 337 the PDE, thus allowing the use of full non-uniform spatial meshes. However, we do 338 339 not apply this change of variables since we use sparse grids to approximate efficiently the solution of (2.2) for higher spatial dimensions. 340

Obviously, if the solution of linear systems in (4.4) turns out too expensive from 341 the computational point of view, the applicability of these schemes remains very 342 limited. However, in the case of the PDE problem here addressed or similar multi-343 dimensional linear problems, due to the simple structure of the matrices \mathcal{A}_i , each 344 linear system of type $(I - \nu \Delta t \mathcal{A}_i)K = G$ can be solved by using $\tilde{L}_i = \prod_{k \neq i}^N M_k$ 345 tridiagonal linear systems of dimension M_i . In order to make easier the reading of 346 this article, the details of this computation are included in Algorithm A.1 in the 347 Appendix. 348

Moreover, another interesting advantage of the AMFR-W-methods (4.4) when applied to the autonomous linear problem (3.4) with $\mathcal{F}(Y) = \mathcal{A}Y, \mathcal{A} = \mathcal{A}_0 + \sum_{i=1}^{N} \mathcal{A}_i$, comes from the fact that the matrix-vector product $\mathcal{F}'(Y_n)K_r^{(N)}$ is simply an extra evaluation of the derivative function $\mathcal{F}(K_r^{(N)})$ and the explicit computation of the matrix \mathcal{A}_0 is not actually necessary.

4.2. θ -method + Gauss-Seidel as W-method. In [33], the authors applied a direct (backwards in time) time-space discretization with finite differences, that can be also interpreted as a W-method (4.1) with fixed time step-size, when a fixed number of iterations of the Gauss-Seidel iterative scheme is used to solve the involvedlinear systems.

More precisely, for a fixed time step-size $\Delta t > 0$, the well-known θ -method applied to (3.4), with $\theta \in [0, 1]$, gives the approximations $W_n \approx Y(t_n), t_n = n\Delta t$, $n = 0, 1, \ldots, M$, by using the formula

362 (4.6)
$$W_{n+1} = W_n + (1-\theta)\Delta t \mathcal{F}(W_n) + \theta \Delta t \mathcal{F}(W_{n+1}).$$

Therefore, when $\theta \neq 0$ and linear problems $\mathcal{F}(Y) = \mathcal{A}Y$ as (3.4) are considered, it is 363 necessary to solve the linear system $(I - \theta \Delta t \mathcal{A})W_{n+1} = \beta_n$, at each time step, with 364 $\beta_n = (I - \theta \Delta t \mathcal{A}) W_n + \Delta t \mathcal{F}(W_n)$. In [33], the Gauss-Seidel iterative linear systems 365 solver is performed until getting an error below a prescribed tolerance. Note that 366 Gauss-Seidel method splits the coefficient matrix $\mathcal{A} = P + R$, where P is the triangular 367 matrix whose entries are the lower-triangular part of \mathcal{A} and its diagonal elements, 368 while R stores its strictly upper-triangular part. By using this splitting, from a 369 starting value $W_{n+1}^{(0)}$, this method computes iteratively approximations $W_{n+1}^{(r)} \approx W_{n+1}$ 370 371 by solving only triangular systems

372 (4.7)
$$(I - \theta \Delta t P) W_{n+1}^{(r)} = \theta \Delta t R W_{n+1}^{(r-1)} + \beta_n, \qquad r = 1, 2, \dots$$

After some algebraic manipulations and taking as a natural choice for the starting value $W_{n+1}^{(0)} = W_n$, these iterations can be written as $W_{n+1}^{(r)} = W_n + \sum_{j=1}^r \hat{K}_j$, $r = 1, 2, \ldots$, where the vectors \hat{K}_j are sequentially computed by

$$(I - \theta \Delta t P)\hat{K}_r = \Delta t \mathcal{A}\left(W_n + \sum_{j=1}^{r-1} \theta \hat{K}_j\right) + \sum_{j=1}^{r-1} (-1)\hat{K}_j \qquad r = 1, 2, \dots$$

If we compare this last formula with (4.1), clearly if this combination of θ -method + 373 Gauss-Seidel iteration (4.6)-(4.7) is performed with a fixed number s of iterations, then 374 the method can be included in the class of W-methods (4.1) with coefficients W = P, 375 $a_{rj} = \theta, \forall j < r, q_{rj} = -1, \forall j < r, b_r = 1, r = 1, \dots, s.$ Therefore, if we include the 376 discretization proposed in [33] in the W-methods framework, we can apply the order 377 conditions given in [22, p.115] or [17, Sec.2.1] (in a similar notation as here), and we 378 obtain that it achieves order 2 (for $s \ge 2$) in time only when $\theta = 1/2$ (Crank-Nicolson 379 scheme), what is in agreement with the results obtained in the aforementioned article 380 [33]. 381

An advantage of expressing the scheme in [33] as a W-method is that it makes 382 easier to compare its computational cost per time step with that of the here proposed 383 **AMFR-W2** method (4.4)-(4.5). The scheme (4.6)-(4.7) with $\theta = 1/2$ and s Gauss-384 Seidel iterations needs to compute one evaluation of the derivative function \mathcal{F} and s 385 triangular linear systems of dimension $L = M_1 \cdots M_N$. On the other hand, **AMFR**-386 **W2** needs to evaluate four times the derivative function (as $\mathcal{F}'(Y_n)K_r^{(N)} = \mathcal{F}(K_r^{(N)})$) 387 and to solve $(2\hat{L}_i)$ tridiagonal linear systems of dimension M_i , per each $i = 1, \ldots, N$. 388 Since the triangular systems of dimension L need approximately $\mathcal{O}(L^2)$ operations, 389 the tridiagonal ones of dimension M_i cost $\mathcal{O}(M_i)$ operations and each evaluation of 390 derivative function involves $\mathcal{O}(L^2)$ operations, we have $\mathcal{O}((s+1)L^2)$ operations for 391 (4.6)-(4.7) and $\mathcal{O}(4L^2 + 2NL)$ operations for **AMFR-W2** (4.4)-(4.5). 392

5. Sparse grids in space. Solving PDE problems as (2.2) on a full tensor prod-393 uct based grid with p^N grid points, with p being the number of grid points in each 394coordinate direction, can become a highly involved computational task, even prohib-395 itive. As the number of underlying forward rates increases, clearly the dimension of 396 the multi-dimensional pricing PDE (2.2) increases as well, so the computational cost 397 of solving the fully discretized problem grows exponentially. Thus, the discretiza-398 tion using this so-called full grid also consumes too much memory. This drawback 399 is referred as the curse of dimensionality. For example, pricing a swaption over five 400 forward rates ruled by the same stochastic volatility, by means of a full grid with 128 401 points per coordinate gives rise to more than one four billion points. The storage 402of such a grid using double precision floating point format will need more than 32 403 404 thousand gigabytes of memory.

Because of the curse of dimensionality, traditional full grid methods, like finite 405differences, finite elements or finite volumes, are not able to price derivatives with 406 high dimensional underlying processes, even in the most powerful supercomputers 407available nowadays. This limitation can be partially overcome by using a family of 408 techniques known as sparse grid methods (see [6], for example). Sparse grids are useful 409410 numerical methods for solving high-dimensional PDEs because they are based on a relatively small number of grid points but also maintain a satisfactory accuracy. More 411 precisely, let d denote the underlying problem's dimensionality and p the number of 412 grid points in one coordinate direction at the boundary. On the one hand, regarding 413 the considered number of degrees of freedom, full grid methods use $O(p^d)$ grid points, 414 while sparse grid discretizations only employ $O\left(p(\log_2 p)^{d-1}\right)$ grid points. On the 415 other hand, concerning accuracy, conventional methods converge at a rate of $O(p^{-2})$ 416 when making use of second order schemes, whereas sparse grid methods converge at the only slightly deteriorated rate of $O(p^{-2} (\log_2 p)^{d-1})$. In [6], Bungartz and Griebel 417 418 present an excellent survey of the fundamentals and the applications of sparse grids, 419with a focus on the solution of PDEs. Sparse grid were introduced in the early 1990s 420 421 for the solution of PDEs by Zenger [44] and Griebel [18].

5.1. Standard sparse grid combination technique. Discretizations on spar-422 se grids require hierarchical data structures. Therefore, specially designed PDE solvers 423 are required, and their implementations become more and more complicated as the 424 dimension of the problem increases [1, 43]. An efficient way to avoid intricate sparse 425 grid implementations is given by the sparse grid combination technique, originally 426 proposed by Griebel, Schneider and Zenger [20]. Basically, the combination technique 427 solves the PDE on several independent and conventional Cartesian smaller-sized grids. 428 Then, the solution in the sparse grid space is approximated by a suitable linear com-429430 bination of these partial solutions on the coarser grid. This solution retains the advertised convergence rate of sparse grid methods if certain error expansions for the 431component approximations exist [7, 8, 3]. Note the rigorous analysis of finite differ-432 ences schemes for the sparse grid combination technique in [39]. Further advantages 433 of the technique are the possibility to solve the problem on each of the constituent 434grids using standard full grid solvers and the inherent parallelism of the method [19]. 435 Let us introduce formally the sparse grid combination technique. We fix a multi-436 index $\boldsymbol{l} = (l_1, l_2, \dots, l_d) \in \mathbb{N}_0^d$ and define its L_1 -norm as $|\boldsymbol{l}|_1 = \sum_{i=1}^d l_i$. In the *d*-dimensional orthohedron $[0, c_1] \times [0, c_2] \times \dots \times [0, c_d]$ $(c_i \in \mathbb{R}_{>0})$, we denote by $\Omega_{\boldsymbol{l}} =$ 437 438 $\Omega_{(l_1,\ldots,l_d)}$ an anisotropic² although full grid having uniform mesh spacing $h_i = 2^{-l_i} c_i$ 439

²Mesh spacing differs in each coordinate direction.

440 in each coordinate direction $i \in \{1, \ldots, d\}$. Let u_l be the conventional finite difference

solution to the PDE on grid Ω_l , extended to $[0, c_1] \times [0, c_2] \times \ldots \times [0, c_d]$ by interpolation.

442 Then, the sparse grid combination solution u_n^s over the sparse grid Ω_n^s with refinement 443 level n is given by the following linear combination

444 (5.1)
$$u_n^s = \sum_{q=0}^{d-1} (-1)^q \cdot \binom{d-1}{q} \cdot \sum_{|\boldsymbol{l}|_1 = n-q} u_{\boldsymbol{l}}.$$

Increasing the level n should give a more accurate solution to the problem. The grid solutions u_1 involved in the inner sum of (5.1) all have $\mathbf{l} = (l_1, \ldots, l_d)$ such that $l_1 + \cdots + l_d = n - q$. The number of elements in each of these grids is $O(2^{n-q})$, regardless of the dimension, and the number of grid solutions in this inner sum is

449
$$\binom{n-q+d-1}{d-1}$$
 and grows like $O((n-q)^{d-1})$. Besides, $\left(\bigcup_{0 \le |\mathbf{l}|_1 < n} \Omega_{\mathbf{l}}\right) \subset \Omega_n^s$. Therefore

the dimension of the sparse grid space on level n is $O(2^n n^{d-1}) = O(h^{-1}(\log_2 h)^{d-1})$ where $h = 2^{-n}$ is the finest grid size. This value can be compared with the full grid space dimension which is $O(2^{nd}) = O(h^{-d})$.

The combination technique works due to the cancellation mechanism of the error terms in the involved grids. This cancellation principle is well known in extrapolation techniques. Indeed, all lower order error terms cancel out in the the combination formula (5.1), see [25] for deep details in dimension two. Thus, the combination technique is able to produce accurate results in reasonable time. Several generalizations of the standard combination technique formula (5.1) have been developed [23].

The combination technique algorithm is embarrassingly parallel since all component grid solutions can be computed in parallel. In general, for refinement level nin d dimensions there are $\sum_{q=0}^{d-1} {\binom{n-q+d-1}{d-1}}$ component grids, which can be solved in parallel. In order to achieve optimal speed-ups one has to carefully deal with load imbalances, even in all those grids at the same refinement level, due to the anisotropic structure of the component grids.

The sparse grid combination technique was initially formulated for elliptic PDEs such as Laplace's and Poisson's equation. Later, it has also been applied to parabolic PDEs, specially for option pricing problems in finance [3, 30, 40, 31, 39, 9, 4, 32, 33, 26, 10, 24, 25]. Here we focus on the implementation of the sparse grid combination technique for parabolic equations. More precisely, we just consider the case where the solution is only needed at the final time, which is frequently the case in finance and particularly in the problem we address.

In this setting, the natural approach is the following. First, solve the parabolic 472 473 equation on each of the full grids involved in the sparse grid combination technique formula (5.1) with a full grid method. Finally, combine these solutions only at the end. 474 This method only requires interpolation from grid values at the final time, but not at 475intermediate time steps. If the numerical error due to the time discretization does not 476 dominate the spatial error, we expect a pointwise rate of convergence proportional to 477 $O(p^{-2}(\log_2 p)^{d-1})$ for our AMFR-W scheme applied to problems with smooth enough 478 initial and boundary data. It is important to notice that interpolation techniques are 479480 required in order to approximate the solution at points not belonging to the sparse grid. The most straightforward approach is to interpolate at those points over all full 481 grids handled by the combination technique, and then add up these results with the 482 appropriate combination technique weights. Note also that the interpolation technique 483has to preserve the order of the used discretization scheme, so that the convergence 484

14

485 result remain valid for the entire domain. Otherwise, the convergence order only holds

486 for grid nodes belonging to all sub-grids and therefore not affected by interpolation. A

tensor based linear interpolation preserves the required order 2 of accuracy for secondorder finite difference discretizations.

5.2. Modified sparse grid combination technique. By means of the previously described standard sparse grid combination technique it is impossible to approximate accurately a Neumann boundary condition for degenerated Cartesian grids having very few points in the corresponding coordinate direction. The approximations of the solution in these grids becomes very poor, thus decreasing the accuracy of the combination technique approximation.

In order to overcome this drawback, a mild modification of the standard sparse grid combination technique (5.1) can be developed, just by forcing a minimum number of discretization steps in all grids involved in the combination procedure. More precisely, all levels in all dimensions start from a small but non zero value ψ , so that the modified combination technique formula reads

500 (5.2)
$$u_n^s = \sum_{q=0}^{d-1} (-1)^q \cdot \binom{d-1}{q} \cdot \sum_{|\boldsymbol{l}|_1 = n-q} u_{\psi \boldsymbol{1}+\boldsymbol{l}},$$

where $\psi \mathbf{1} + \mathbf{l} = (\psi + l_1, \dots, \psi + l_d)$. This modified sparse grid combination technique 501working over a modified sparse grid $\Omega_n^{s,\psi}$ produces more accurate approximations 502 [4] at the cost of increasing the consumed time and memory. Although the number 503of subproblems to be solved is exactly the same as before, the number of degrees 504of freedom associated to each subproblem increases. In fact, the number of grid 505points in the combined sparse grid increases from $O(2^n n^{d-1})$ in the standard one 506 to $O(2^{n+d\psi}n^{d-1})$ in the modified one. In this new setting ψ should be kept small 507 (specifically $\psi = 1$ or 2 in the present work), otherwise the new modified combination 508 technique will suffer soon the curse of dimensionality. 509

6. Numerical results. In this Section we present the obtained numerical results when the previously described methodologies are applied. More precisely, we show and discuss the results obtained by using the AMFR-W method with full grid, standard and modified sparse grids combination techniques to conveniently cope with the proposed homogeneous Neumann boundary conditions in the particular case $\beta = 1$.

For all products we will use the data presented in Table 1 where we consider the tenor structure $0 = T_0 < 0.5 < 1.0 < \cdots < 2.5 < 3 = T_{10}$ in years, with constant periods $\tau = T_{i+1} - T_i = 0.5$.

i	T_i	$F_i(0)$	$lpha_i$	i	T_i	$F_i(0)$	$lpha_i$
0	0	0.0112	0	3	1.5	0.0126	0.2221
1	0.5	0.0118	0.2366	4	2	0.0130	0.2068
2	1	0.0122	0.2145	5	2.5	0.0135	0.1932
			Тав	LE 1			

Hypothetical market data (LIBOR rates and volatilities) used in pricing. Strike rate K = 0.011.

The spatial domain is defined by $F^{max} = 0.04$ and $V^{max} = 3.5$, thus upper boundaries were settled between 3 and 4 times the point of interest at which we evaluate the pricing of the interest rate derivative. In the cases where the analytical solution is not available, we first compute reference solutions using the proposed space and time discretizations over classical full grids. These solutions will serve to assess on the accuracy of the proposed sparse grids methods in space. The designed algorithms were implemented using C++ (GNU C++ compiler 8.3.1) and double precision. Besides, all numerical experiments have been performed in a machine with 16GBytes of RAM and four multicore Intel Xeon CPUs E5-2620 v4 clocked at 2.10GHz, each one with eight cores.

6.1. Numerical results with full grids. The first test to validate the proposed numerical methodologies consists of pricing a caplet without considering stochastic volatility, that is to say, under the classical LMM. This test is a sanity check, since the analytical pricing formula is known for caplets, the so-called Black-Scholes's formula for caplets ([5, equation 1.26]). More precisely, we start pricing the caplet with maturity T_1 and payoff $\tau_1(F_1(T_1) - K)^+$ paid at time T_2 , under the data of Table 1, with strike rate K = 0.011. The present intrinsic value of the caplet is given by $P(0, T_2)\tau_1(F_1(T_1) - K)^+$. The exact price of this product given by Black-Scholes' formula is 6.058877 basis points (bps, 1 bp $= 10^{-4}$).

In order to price this caplet using the PDE approach presented here it is con-537538 venient to consider the terminal probability measure associated with choosing the bond $P(0,T_2)$ as numeraire. Thus, the price of this product is given by the solution of the PDE (2.2) (with $\sigma = 0$) multiplied by $P(0, T_2) = \frac{1}{1 + \tau F_0(0)} \frac{1}{1 + \tau F_1(0)}$, i.e. 540 $P(0,T_2)u(0,F_1,V)$. Once obtained the PDE solution on the last time slice, the price 541 of the caplet is obtained through interpolation in space, by means of multilinear in-542terpolation, thus maintaining order two in space. In Table 2 full grid solutions are presented for levels from 6 to 13 in space and considering 4, 8, 16 and 256 time steps. 544The interpolation in space for the last time slice was done in $F_1 = 0.0118$ and V = 1. 545The column labelled as Solution shows the PDE solution in bps, and the column for 546 the error measures the absolute distance of the numerical solution to the exact one, in bps as well. The execution time was measured in seconds in all the experiments in this 548 work. The "grid points" column displays the number of grid points employed in the 549full space meshes at each time discretization. Since the method is order three in time, 550few time steps could be considered in real pricing applications. For the space level 13 552 (for the forward and the volatility) and when using 256 time steps, the method was able to recover the exact solution up to the 8-th decimal digit. Nevertheless, in this 553 case the full grid method required almost three hours, all space meshes in all time 554555slices with more than 67 million points.

556 Once we have checked the correct behaviour of the full grid method, which will be 557 used in the sparse grid combination technique, we compute full grid reference solutions 558 for financial products without exact prices. They will be used in order to assess 559 the correctness of the upcoming sparse grid combination technique implementation. 560 Therefore, in order to minimize errors due to the time discretization, 256 time steps 561 will be chosen for the rest of the full grid tests in this section. In Table 3 the computed 562 prices of the previous caplet under the stochastic volatility framework are shown.

Next, we deal with the pricing of $T_a \times (T_b - T_a)$ European swaptions. In Table 4, first the results for the 0.5×1 swaption are given. Note that under this full grid framework it is not possible to price this product in reasonable computational times past refinement level 9, due to the high number of involved spatial grid points. Then, the results for the 0.5×1.5 swaption are also shown. Once more, full grid pricing is only achievable on the lower grid levels.

569 6.2. Numerical results with the standard sparse grid combination tech-570 nique. In this section, by means of the standard sparse grid combination technique, 571 we price not only the previous caplets and swaptions, but also swaptions involving

		4 time steps			8 time steps		
Level	Solution	Error	Time	Solution	Error	Time	Grid points
6	6.084214	2.533666×10^{-2}	0.01	6.082708	2.383026×10^{-2}	0.02	4225
7	6.065957	7.079711×10^{-3}	0.05	6.064274	5.396418×10^{-3}	0.09	16641
8	6.063526	4.648521×10^{-3}	0.15	6.061980	3.102002×10^{-3}	0.27	66049
9	6.060148	1.270262×10^{-3}	0.55	6.058939	6.133732×10^{-5}	0.96	263169
10	6.060300	1.422802×10^{-3}	2.14	6.059077	1.992885×10^{-4}	4.23	1050625
11	6.060237	1.359825×10^{-3}	9.39	6.059021	1.431228×10^{-4}	18.69	4198401
12	6.060236	1.358165×10^{-3}	40.47	6.059019	1.416918×10^{-4}	80.45	16785409
13	6.060226	1.348640×10^{-3}	165.64	6.059011	1.331767×10^{-4}	328.15	67125249
		16 time steps			256 time steps		
Level	Solution	Error	Time	Solution	Error	Time	Grid points
6	6.082540	2.366282×10^{-2}	0.04	6.082513	2.363549×10^{-2}	0.47	4225
7	6.064109	5.231082×10^{-3}	0.16	6.064079	5.201370×10^{-3}	1.89	16641
8	6.061870	2.992441×10^{-3}	0.49	6.061840	2.962119×10^{-3}	7.28	66049
9	6.058832	4.556637×10^{-5}	1.97	6.058802	7.582048×10^{-5}	30.77	263169
10	6.058975	9.701419×10^{-5}	8.33	6.058944	6.680796×10^{-5}	131.75	1050625
11	6.058919	4.157937×10^{-5}	37.19	6.058889	1.139059×10^{-5}	593.13	4198401
12	6.058918	4.041366×10^{-5}	163.01	6.058888	1.022896×10^{-5}	2558.92	16785409
13	6.058910	3.191855×10^{-5}	655.84	6.058879	1.735259×10^{-6}	10415.57	67125249
			r	TABLE 2			

Full grid method, caplet with expiry T_1 , $\sigma = 0$. Prices and errors are shown in bps.

Level	Solution	Time	Grid points	Level	Solution	Time	Grid points
6	6.050103	0.48	4225	10	6.023799	135.29	1050625
7	6.029510	1.93	16641	11	6.023737	597.82	4198401
8	6.026929	7.51	66049	12	6.023734	2557.62	16785409
9	6.023665	31.06	263169	13	6.023725	10505.02	67125249
			TA	BLE 3			

Full grid method, caplet with expiry T_1 , $\sigma = 0.3$, $\phi_1 = 0.4$, 256 time steps.

-	Level	Solution	Time	Grid points	Level	Solution	Time	Grid points					
X	6	13.002003	71.33	274625	8	12.981320	5111.14	16974593					
0.5	7	12.984709	590.97	2146689	9	12.980459	43325.53	135005697					
1.5	Level	Solution	Time	Grid points	Level	Solution	Time	Grid points					
×	3	23.952705	1.70	6561	5	21.577149	474.93	1185921					
0.5	4	21.765442	28.45	83521	6	21.486079	8122.91	17850625					

Full grid method, 0.5×1 and 0.5×1.5 swaptions, $\sigma = 0.3$, $\phi_i = 0.4$, i = 1, 2, 3, 256 time steps.

more underlying forward interest rates, thus dealing with a high dimensional setting. 572573 As usual, we are also interested in the values of these derivatives at the last time cut for the values of the forward rates depicted in Table 1 and V = 1, which define the 574spatial point where the value of the solution of the PDE is computed. In order to 575obtain the solution given by the sparse grid combination technique at this point, the 576numerical solution on each grid involved in the combination technique is interpolated 577 at this point with multilinear interpolation. Next, all these values are introduced in 578the combination technique formula (5.1), thus obtaining the price provided by the 579standard sparse grid combination technique. 580

Moreover, sparse grid combination techniques have been implemented to take 581advantage of shared memory parallel computers. The code was optimized and paral-582lelized using OpenMP framework [45], version 4.5. In order to deal with the previously 583 584 mentioned load imbalances it is crucial to use a dynamic schedule to assign the involved full grids to threads. In this way, OpenMP assigns one grid to each thread. 585When the thread finishes, it will be assigned the next mesh that has not been as-586signed yet. The speedup of the parallelized version is almost equal to the number of 587 588 available computing cores, in our case 32. This optimal speedup is due to the fact that communication between processors only takes place at the final step in order to concentrate the solutions over all grids to a single scalar value.

In Table 5 we price the caplet with maturity T_1 under the framework without stochastic volatility, whose exact price was 6.058877 basis points. The results in this 592 Table are to be compared with those of Table 2. The accuracy of each solution is only slightly worse in this case, although the computing time is much lower due the 594much less number of involved grid points. For example, with the full grid approach, 595the solution using 256 time steps and refinement level in space 13 employed 10415.57 596seconds to achieve an error 1.735259×10^{-6} in basis points, while the standard sparse 597 grid combination technique just needed 569.26 seconds to attain almost the same 598 accuracy, an error of 1.047159×10^{-5} . The reduction in the number of employed grid 599 600 points is also shown in Table 2.

		4 time steps			8 time steps		
Level	Solution	Error	Time	Solution	Error	Time	Grid points
6	6.063081	4.202988×10^{-3}	0.02	6.053883	4.994378×10^{-3}	0.02	385
7	6.120850	6.197245×10^{-2}	0.02	6.110998	5.212012×10^{-2}	0.02	833
8	6.067165	8.287799×10^{-3}	0.02	6.059333	4.557323×10^{-4}	0.03	1793
9	6.061071	2.192917×10^{-3}	0.03	6.056314	2.563871×10^{-3}	0.05	3841
10	6.061681	2.803747×10^{-3}	0.07	6.059090	2.119811×10^{-4}	0.12	8193
11	6.059735	8.569127×10^{-4}	0.20	6.058133	7.450762×10^{-4}	0.35	17409
12	6.060156	1.278608×10^{-3}	0.58	6.058787	9.054229×10^{-5}	1.21	36865
13	6.060211	1.333075×10^{-3}	2.43	6.058951	7.311969×10^{-5}	4.53	77825
14	6.060223	1.345681×10^{-3}	9.00	6.058998	1.206373×10^{-4}	17.90	163841
		16 time steps			256 time steps		
Level	Solution	Error	Time	Solution	Error	Time	Grid points
6	6.052778	6.099826×10^{-3}	0.02	6.052738	6.139440×10^{-3}	0.05	385
7	6.110144	5.126590×10^{-2}	0.03	6.110118	5.124050×10^{-2}	0.09	833
8	6.059007	1.291280×10^{-4}	0.04	6.058984	1.066738×10^{-4}	0.22	1793
9	6.056190	2.687266×10^{-3}	0.07	6.056164	2.713620×10^{-3}	0.71	3841
10							
	6.059029	1.514420×10^{-4}	0.20	6.058998	1.205596×10^{-4}	2.31	8193
11	6.059029 6.058034	$\frac{1.514420 \times 10^{-4}}{8.433494 \times 10^{-4}}$	0.20	6.058998 6.058003	$\frac{1.205596 \times 10^{-4}}{8.741697 \times 10^{-4}}$	2.31 8.47	8193 17409
11 12	6.059029 6.058034 6.058698	$\frac{1.514420 \times 10^{-4}}{8.433494 \times 10^{-4}}$ $\frac{1.800588 \times 10^{-4}}{1.800588 \times 10^{-4}}$	0.20 0.64 2.33	6.058998 6.058003 6.058667	$\frac{1.205596 \times 10^{-4}}{8.741697 \times 10^{-4}}$ $\frac{2.108407 \times 10^{-4}}{10^{-4}}$	2.31 8.47 34.02	8193 17409 36865
$ \begin{array}{c} 11\\ 12\\ 13\\ \end{array} $	6.059029 6.058034 6.058698 6.058852	$\frac{1.514420 \times 10^{-4}}{8.433494 \times 10^{-4}}$ $\frac{1.800588 \times 10^{-4}}{2.520960 \times 10^{-5}}$	0.20 0.64 2.33 8.99	6.058998 6.058003 6.058667 6.058822	$\frac{1.205596 \times 10^{-4}}{8.741697 \times 10^{-4}}$ $\frac{2.108407 \times 10^{-4}}{5.565743 \times 10^{-5}}$	$ \begin{array}{r} 2.31 \\ 8.47 \\ 34.02 \\ 141.73 \end{array} $	8193 17409 36865 77825
11 12 13 14	6.059029 6.058034 6.058698 6.058852 6.058897	$\begin{array}{c} 1.514420 \times 10^{-4} \\ 8.433494 \times 10^{-4} \\ 1.800588 \times 10^{-4} \\ 2.520960 \times 10^{-5} \\ 1.981956 \times 10^{-5} \end{array}$	0.20 0.64 2.33 8.99 35.69	$\begin{array}{r} 6.058998\\ \hline 6.058003\\ \hline 6.058667\\ \hline 6.058822\\ \hline 6.058867\\ \end{array}$	$\begin{array}{c} 1.205596\times10^{-4}\\ 8.741697\times10^{-4}\\ 2.108407\times10^{-4}\\ 5.565743\times10^{-5}\\ 1.047159\times10^{-5} \end{array}$	$ \begin{array}{r} 2.31 \\ 8.47 \\ 34.02 \\ 141.73 \\ 569.26 \\ \end{array} $	8193 17409 36865 77825 163841

Sparse grid combination technique, caplet with expiry T_1 , $\sigma = 0$.

Next, in Table 6 the results for the previous caplet under the stochastic volatility 601 framework are shown. These results are to be compared with those of Table 3. Then, 602 Tables 7 and 8 show the prices given by the standard sparse grid combination tech-603 nique for 0.5×1 and 0.5×1.5 swaptions under stochastic volatility. These results are 604 to be compared with those of Table 4. Clearly, the standard sparse grid combination 605technique outperforms the full grid approach. Besides, the sparse method is able to 606 cope with higher resolution levels, thus allowing to price successfully the 0.5×1.5 607 swaption. Note that this was not possible with the full grid approach, see Table 4. 608

Level	Solution	Time	Level	Solution	Time
6	6.057668	0.06	11	6.022082	8.40
7	6.095685	0.10	12	6.023257	33.96
8	6.025848	0.25	13	6.023595	141.41
9	6.018222	0.70	14	6.023693	569.07
10	6.021834	2.32			
		TAI	SLE 6		

Sparse grid combination technique, caplet with expiry T_1 , $\sigma = 0.3$, $\phi_1 = 0.4$, 256 time steps.

Level	Solution	Time	Grid points	Level	Solution	Time	Grid points
8	12.311172	0.67	8705	12	13.205324	89.80	219137
9	13.024747	2.00	19713	13	12.993536	360.72	483329
10	13.616333	6.77	44289	14	12.971783	1399.46	1060865
11	13.525821	24.42	98817	15	12.973900	5755.68	2318337
			Та	BLE 7			

Sparse grid combination technique, 0.5×1 swaption, $\sigma = 0.3$, $\phi_1 = \phi_2 = 0.4$, 256 time steps.

	8 time	steps	256 tin	ne steps	
Level	Solution	Time	Solution	Time	Grid points
12	21.935448	8.50	21.936574	271.54	1064961
13	21.842901	31.75	21.844522	998.51	2439169
14	21.609183	119.97	21.610055	3809.31	5550081
15	21.707363	461.28	21.708001	14866.96	12554241
16	21.519917	1838.36	21.516402	59010.15	28246017
17	21.483062	7315.95	21.478930	235912.03	63242241
	•		TABLE 8		

Sparse grid combination technique, 0.5×1.5 swaption $\sigma = 0.3$, $\phi_1 = \phi_2 = \phi_3 = 0.4$.

Finally, in Tables 9 and 10, 0.5×2 and 0.5×2.5 swaptions are priced under 609 stochastic volatility. The curse of dimensionality makes impossible to price these 610 products with full grid approaches. In order to speedup the convergence of the sparse 611 grid method, a useful technique is to consider a computational domain such that the 612 613 point of interest is in the neighbourhood of the center of the domain. This strategy easily improves sparse grid results. In fact, in that region is where the sparse grid 614 contains more points. Indeed, the central point belongs to all non degenerated grids 615 involved in the standard sparse grid combination technique. The improvement in 616 accuracy can be observed in Table 10, where the upper boundaries of the forward 617 rates, F^{max} , were shrunk from 0.04 to 0.02. 618

	4 time	e steps	8 time steps		
Level	Solution	Time	Solution	Time	
14	35.341806	180.34	35.346408	360.57	
15	34.388334	669.85	34.425087	1335.40	
16	32.115380	2561.87	32.122101	5133.81	
17	30.639336	10058.18	30.641664	20076.39	
18	30.881086	40097.11	30.918448	80268.57	
19	30.822037 239746.17		30.797087	479681.87	
		TABLE 9			

Sparse grid combination technique, 0.5×2 swaption, $\sigma = 0.3$, $\phi_1 = \ldots = \phi_4 = 0.4$.

In order to price interest rate derivatives involving more underlying forward rates using this approach, the proposed algorithm should be implemented to run on a cluster of processors (distributed memory machines). Since the communications between processors is minimal, the technique scales optimally. This extra layer of parallelism would bring also a further reduction on the previous execution times, thus allowing to stress the method with higher resolution levels.

625 **6.3.** Numerical results with the modified sparse grid combination tech-626 nique. Our last set of numerical experiments aims at showing that the modified 627 sparse grid technique defined by (5.2) is able to improve the performance (accuracy 628 and computing time) of the standard sparse grid combination technique given by 629 expression (5.1), specially in moderately high dimensions.

	$F^{max} = 0.04$										
	2 time	e steps	4 time	e steps	8 time steps						
Level	Solution	Time	Solution	Time	Solution	Time					
16	54.923235	3812.17	53.634032	7565.44	53.512820	15084.20					
17	39.023559	14245.68	41.296346	28565.06	41.328455	56819.84					
18	39.780626	55108.12	38.377139	109843.14	38.823376	220825.17					
19	41.230567	285159.96	41.970631	570320.92	41.599974	1140639.84					
	$F^{max} = 0.02$										
	2 time	e steps	4 time	e steps	8 time steps						
Level	Solution	Time	Solution	Time	Solution	Time					
16	42.250960	3812.17	43.016757	7565.44	42.777426	15084.20					
17	42.116312	14245.68	41.406882	28565.06	41.842625	56819.84					
18	42.991274	55108.12	42.746717	109843.14	42.737987	220825.17					
19	42.446354	285159.96	42.892002	570320.92	42.838119	1140639.84					
			TABLE 1	0							

Sparse grid combination technique, 0.5×2.5 swaption, $\sigma = 0.3$, $\phi_1 = \ldots = \phi_5 = 0.4$.

As in the previous cases, we start with the sanity test of the pricing of the caplet 630 with expiry T_1 under the classical LMM. Table 11 gathers the behaviour of the mod-631 ified technique when pricing this caplet. Firstly, we compare Table 11 with Table 632 5 originated with the standard combination technique. With $\psi = 1$, the modified 633 technique is able to obtain an accuracy of 1.047853×10^{-5} with level equal 12 in less 634 than a hundred seconds. In contrast, the standard sparse grid technique required a 635 higher refinement level of 14 and employed more than five hundred seconds to obtain 636 a similar accuracy. Also note that with $\psi = 2$ and the refinement level 10, the mod-637 ified combination technique is able to get better results, an error of 7.625043×10^{-6} 638 639 in just over 25 seconds. Moreover, while the obtained order of convergence in space for the standard combination technique is slightly worse than two, with this modified 640 = 2. The comparison with full grid method results method is almost two when 641 shown in Table 2 could be summarized by noting that with $\psi = 2$ the modified sparse 642 grid technique is able to obtain an error less than 1.735259×10^{-6} in less than five 643 hundred seconds, while the full grid approach needed almost 2.9 hours. 644

		$\psi = 1$			$\psi = 2$				
n	Solution	Error	Time	#points	Solution	Error	Time	#points	
7	6.056324	2.55×10^{-3}	0.20	2817	6.057952	9.25×10^{-4}	0.70	10241	
8	6.058978	1.00×10^{-4}	0.44	6145	6.058684	1.93×10^{-4}	2.04	22529	
9	6.058005	8.72×10^{-4}	1.61	13313	6.058822	5.51×10^{-5}	7.08	49153	
10	6.058666	2.11×10^{-4}	6.35	28673	6.058870	7.62×10^{-6}	26.06	106497	
11	6.058822	5.52×10^{-5}	24.62	61441	6.058874	4.08×10^{-6}	105.78	229377	
12	6.058867	1.04×10^{-5}	96.95	131073	6.058877	9.15×10^{-7}	478.10	491521	
13	6.058873	4.45×10^{-6}	473.16	278529	6.058877	3.17×10^{-7}	1991.75	1048577	
14	6.058877	1.07×10^{-6}	1968.99	589825	6.058878	9.91×10^{-8}	7868.34	2228225	

Table 11

Modified sparse grid combination technique, caplet with maturity T_1 , $\sigma = 0$, 256 time steps, $F^{max} = 0.04$.

Finally, Tables 12 and 13 show the results for 4 and 6 dimensional PDEs in space, respectively. These Tables are to be compared with the corresponding Tables and 10 generated with the standard combination technique. We observe that the higher the dimensionality of the problem the lower ψ should be, otherwise the curse of dimensionality will appear soon again. Table 12 shows that with $\psi = 1$ the modified method is able to obtain in just 15 minutes a similar accuracy to the one obtained with the standard approach in more than 246 minutes. With $\psi = 2$ we observe that 5 decimal digits are stabilized in the modified combination technique. Regarding Table 13, with $\psi = 1$ three decimal digits are stabilized already in level 13. Also note that for the 6 dimensional PDEs in space reported in Table 13, in our machine the modified method is not able to go further level 12 with $\psi = 2$ in a reasonable computational time due to the curse of dimensionality. Nevertheless, the accuracy recovered for level 12 is remarkable. Finally, notice that adding points entails better performance than increasing the level of the sparse grid.

		$\psi = 1$			$\psi = 2$	
Level	Solution	Time	Grid points	Solution	Time	Grid points
9	21.511031	6.11	114689	21.519347	107.59	763905
10	21.297979	19.02	262145	21.491243	328.03	1765377
11	21.600032	60.94	593921	21.485699	1051.35	4038657
12	21.595254	209.27	1335297	21.466814	3807.11	9158657
13	21.472738	806.72	2981889	21.464192	13921.97	20611073
			TABLE	12		

Modified sparse grids combination technique, swaption 0.5×1.5 , $\sigma = 0.3$, $\phi_i = \phi_2 = \phi_3 = 0.4$, 16 time steps, $F^{max} = 0.04$.

	$\psi = 1$		$\psi = 2$	
Level	Solution	Time	Solution	Time
12	42.624837	3255.48	42.723635	420696.70
13	42.828046	10476.64	_	_
14	42.686665	34859.98	-	_
15	42.702808	121471.65	_	_
TABLE 13				

Modified sparse grid combination technique, swaption 0.5×2.5 , $\sigma = 0.3$, $\phi_1 = \ldots = \phi_5 = 0.4$, 4 time steps, $F^{max} = 0.02$.

659 7. Conclusions. In this work we have mainly developed a new numerical methodology which combines high order time discretization algorithms with a sparse 660 grids modified combination technique to solve high dimensional PDE problems aris-661 ing in finance. More precisely, we have focused on the numerical solution of the PDE 662 formulation proposed in [33] for pricing a large variety of interest rate derivatives, 663 when the underlying forward rates follow a SABR-LMM model. For this purpose, 664665 we have proposed the use of high order in time AMFR-W methods, thus allowing the use of larger time steps. Moreover, a suitable splitting of the involved operators 666 additionally contributes to the computational time reduction for a given accuracy. 667 As the PDE problem becomes high dimensional in space when the particular interest 668 669 rate derivative requires the consideration of a large number of forward rates (each one giving rise to one spatial dimension), the application of AMFR-W methods on 670 sparse grids with combination technique turns out to be very efficient to obtain the 671 pricing in reasonable computational times. As illustrated in the section of numeri-672 cal results, parallel implementations of the algorithms based on OpenMP framework 673 674 lead to a significant speed up of the computations. As indicated, an appropriate load imbalances management provides an optimal speed up, which is almost equal to the 675 676 number of available computer cores. All computer implementations have been carried out from scratch. Another relevant innovative aspect comes from the suitable con-677 sideration of new homogeneous Neumann boundary conditions, instead of Dirichlet 678 ones in [33]. This consideration avoids the numerical difficulties associated to the 679 680 presence of boundary layers in the outflow boundaries when the parameter β is not 281 zero, specially in the advection dominated regime. Moreover, they motivate the introduction of a modified combination technique to cope with a certain decrease in the accuracy of the standard combination technique, which mainly comes from the inaccuracy of approximations obtained with some degenerated grids included in the sparse grids combination expression. Numerical results also illustrate the advantages of the proposed modified combination technique with respect to the standard version.

Although this article focuses on the PDE formulation of the very relevant financial 687 problem of pricing interest rate derivatives, the proposed methodology can be applied 688 to a large variety of models involving high dimension PDE formulations which arise 689 not only in finance but also in other disciplines in sciences and engineering. For 690 example, in finance high dimension PDE problems related to the pricing of basket 691 692 options or the computation of the XVA associated to portfolios could be considered. In computational biology, the same happens with problems related to gene networks 693 or synthetic biology. 694

695 Appendix A. Appendix.

Proof of Lemma 3.1: Since $M_i - m_i \ge 0$, $j_i - m_i \ge 0$, for all i, it is clear that $\vartheta(m_1, m_2, \ldots, m_N) \le \vartheta(\mathbf{j}) \le \vartheta(M_1, M_2, \ldots, M_N)$, for all $\mathbf{j} \in \mathcal{I}_N$. It is easy to see that $\vartheta(m_1, m_2, \ldots, m_N) = m_1$ and $\vartheta(M_1, M_2, \ldots, M_N) = M_T + m_1 - 1$. 699 On the other hand, calling $Q_l := M_l - m_l \ge 0$, if $\vartheta(\mathbf{j}) = \vartheta(\mathbf{k})$ for $\mathbf{j}, \mathbf{k} \in \mathcal{I}_N$,

700 (A.1)
$$q_1 = -\sum_{l=2}^{N} \left(q_l \prod_{r=1}^{l-1} (Q_r + 1) \right), \qquad q_l := j_l - k_l, \, \forall 1 \le l \le N,$$

where the differences q_l are integers that satisfy $|q_l| \leq Q_l$, $1 \leq l \leq N$. Let us suppose that N > 2 (when N = 2 is much simpler). From (A.1), we obtain that

703 (A.2)
$$q_1 = -\left(q_2 + \sum_{l=3}^{N} \left(q_l \prod_{r=2}^{l-1} (Q_r + 1)\right)\right) (Q_1 + 1)$$

Therefore, q_1 is a multiple of the positive integer $(Q_1 + 1)$. Moreover, as $-Q_1 \leq q_1 \leq Q_1$, so necessarily q_1 must be zero. Therefore, because of (A.2), we obtain

$$q_2 = -\sum_{l=3}^{N} \left(q_l \prod_{r=2}^{l-1} (Q_r + 1) \right),$$

which is the same formula as in (A.1), although starting from r = 2 instead of r = 1. Applying a similar procedure as for q_1 , we get that necessarily $q_2 = 0$. Inductively, we obtain that $q_l = 0, i = 1, 2, ..., N$, so $\mathbf{j} = \mathbf{k}$ and the map ϑ is injective. Since clearly the two sets \mathcal{I}_N and $\{m_1, m_1 + 1, ..., M_T + m_1 - 1\}$ have the same number M_T of elements, ϑ is a bijection.

As a consequence of Lemma 3.1, for all $J \in \{m_1, m_1 + 1, \ldots, M_T + m_1 - 1\}$, there exists a unique multi-index $\mathbf{j} \in \mathcal{I}_N$ given by $\mathbf{j} = \vartheta^{-1}(J)$. In practice, it is necessary to compute this inverse when we manipulate finite differences. An efficient way to calculate it is to use the modulo operation, i.e., $(a \mod n)$ is the remainder of the Euclidean division of a by n.

T14 LEMMA A.1. For every integer $J \in \{m_1, m_1+1, \ldots, M_T+m_1-1\}$, the components of the unique multi-index $\mathbf{j} = (j_1, \ldots, j_N) = \vartheta^{-1}(J) \in \mathcal{I}_N$ can be written as $j_i =$ 716 $m_i + b_i, i = 1, ..., N$ where the integers b_i satisfy

(A.3)
$$c_1 = J - m_1, \qquad b_1 = c_1 \mod (M_1 - m_1 + 1),$$
$$c_i = \frac{c_{i-1} - b_{i-1}}{M_{i-1} - m_{i-1} + 1}, \qquad b_i = c_i \mod (M_i - m_i + 1), \quad i = 2, \dots, N.$$

Algorithm A.1 Procedure to solve linear systems of type $(I - \nu \Delta t A_i)K = G$ Define a matrix Q of dimension M_i ; K = G

$$\begin{array}{l} \mbox{if } i = 1 \ \mbox{the large line is the large line$$

Proof of Lemma A.1: With the same notation as in the proof of Lemma 3.1, $b_l = j_l - m_l \in \{0, 1, \dots, Q_l\}$ for all $l = 1, \dots, N$, equation (3.2) turns into

$$c_1 := J - m_1 = b_1 + \sum_{l=2}^N b_l \prod_{r=1}^{l-1} (Q_r + 1) = b_1 + (Q_1 + 1) \left(b_2 + \sum_{l=3}^N b_l \prod_{r=2}^{l-1} (Q_r + 1) \right),$$

then $b_1 = c_1 \mod (Q_1 + 1) \in \{0, 1, \dots, Q_1\}, j_1 = b_1 + m_1$, and

$$c_2 := \frac{c_1 - b_1}{Q_1 + 1} = b_2 + \sum_{l=3}^N b_l \prod_{r=2}^{l-1} (Q_r + 1).$$

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Clearly, we can apply again the mod operation to c_2 , obtaining $b_2 = c_2 \mod (Q_2 + 1) \in \{0, 1, \dots, Q_2\}$ and $j_2 = b_2 + m_2$. The proof is completed by iteration.

⁷²⁰ In order to help the reader interested in computing every directional linear system

of type $(I - \nu \Delta t A_i)K = G$ of dimension $M, i = 1, \dots, N$, the procedure to solve them

 $_{722}$ is presented in Algorithm A.1, when the *i*-direction and the right-hand side vector

- 723 G are given. We must observe that a new bijection ϑ_{N-1} of type (3.2) is used there
- for multi-indices of dimension N-1 with $m_i = 1, M_i = \hat{L}_{i+1}, i = 1, \dots, N-1$.
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